# Spectral analysis of two doubly infinite Jacobi operators 

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## Spectral Theory and Applications

conference in memory of Boris Pavlov

$$
\text { March 13, } 2016
$$

## (1) Introduction

## (2) Spectral resolution of A

(3) Spectral resolution of $B$

## Two doubly-infinite Jacobi matrices

We analyze the spectral properties of Jacobi operators $A$ and $B$ acting on vectors of the standard basis of $\ell^{2}(\mathbb{Z})$ as:

$$
A e_{n}=q^{-n+1} e_{n-1}+q^{-n} e_{n+1}, \quad n \in \mathbb{Z}
$$

and

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B e_{n}=e_{n-1}+\alpha q^{-n} e_{n}+e_{n+1}, \quad n \in \mathbb{Z},
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The spectrum of any associated semi-infinite Jacobi operator is never known explicitly $(\alpha \neq 0)$ but is expressible in terms of zeros of certain special functions.functions

## Special functions $1 / 2$ - basic hypergeometric series

- Let $0<q<1, r, s \in \mathbb{Z}_{+}$. Recall the basic hypergeometric function

$$
r \phi_{s}\left[\begin{array}{llll}
a_{1}, & a_{2}, & \ldots & a_{r} \\
b_{1}, & b_{2}, & \ldots & b_{s}
\end{array} ; q, z\right]
$$

is defined by the power series

$$
\sum_{n=0}^{\infty} \frac{\left(a_{1} ; q\right)_{n}\left(a_{2} ; q\right)_{n} \ldots\left(a_{r} ; q\right)_{n}}{\left(b_{1} ; q\right)_{n}\left(b_{2} ; q\right)_{n} \ldots\left(b_{s} ; q\right)_{n}} \frac{(-1)^{(s-r+1) n} q^{(s-r+1) n(n-1) / 2}}{(q ; q)_{n}} z^{n}
$$

where $z, a_{1}, a_{2}, \ldots, a_{r} \in \mathbb{C}, b_{1}, b_{2}, \ldots, b_{s} \in \mathbb{C} \backslash q^{\mathbb{Z}_{-}}$and

$$
(a ; q)_{0}=1, \quad(a ; q)_{n}=(1-a)(1-a q) \ldots\left(1-a q^{n-1}\right)
$$

is the $q$-Pochhammer symbol.

## Special functions 1/2-basic hypergeometric series

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- Here we will need only ${ }_{0} \phi_{1}$ and ${ }_{1} \phi_{1}$.


## Special functions 2/2-theta functions

- The theta function:

$$
\theta_{q}(z):=(z ; q)_{\infty}(q / z ; q)_{\infty}=\frac{1}{(q ; q)_{\infty}} \sum_{n=-\infty}^{\infty} q^{n(n-1) / 2}(-z)^{n}
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- Jacobi's theta functions:

$$
\begin{aligned}
& \vartheta_{1}(z \mid q)=\mathrm{i} q^{1 / 4} e^{-\mathrm{i} z}\left(q^{2} ; q^{2}\right)_{\infty} \theta_{q^{2}}\left(e^{2 \mathrm{i} z}\right) \\
& \vartheta_{2}(z \mid q)=q^{1 / 4} e^{-\mathrm{i} z}\left(q^{2} ; q^{2}\right)_{\infty} \theta_{q^{2}}\left(-e^{2 \mathrm{i} z}\right) \\
& \vartheta_{3}(z \mid q)=\left(q^{2} ; q^{2}\right)_{\infty} \theta_{q^{2}}\left(-q e^{2 \mathrm{i} z}\right) \\
& \vartheta_{4}(z \mid q)=\left(q^{2} ; q^{2}\right)_{\infty} \theta_{q^{2}}\left(q e^{2 \mathrm{i} z}\right)
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## Operators associated with the Jacobi matrix $\mathcal{A}$

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- Let $D:=\left\{\psi \in \ell^{2} \mid \mathcal{A} \psi \in \ell^{2}\right\}$. By using the theory of self-adjoint extensions and simple structure of matrix $\mathcal{A}$ one gets:


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## Proposition (self-adjoint extensions)

For $t \in \mathbb{R} \cup\{\infty\}$, operators $A_{t}$, acting as $A_{t} \psi=\mathcal{A} \psi$, with domains

$$
\operatorname{Dom} A_{t}=\left\{\psi \in D \mid \lim _{n \rightarrow \infty} q^{-n}\left(\psi_{2 n+1}+t \psi_{2 n}\right)=0 \wedge \lim _{n \rightarrow \infty} q^{-n}\left(q \psi_{2 n-1}-t \psi_{2 n}\right)=0\right\}
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if $t \in \mathbb{R}$, or

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\operatorname{Dom} A_{\infty}=\left\{\psi \in D \mid \lim _{n \rightarrow \infty} q^{-n} \psi_{2 n}=0\right\}
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are all self-adjoint extensions of $A$.

- In addition,

$$
\sigma_{c}\left(A_{t}\right)=\sigma_{\operatorname{ess}}\left(A_{t}\right)=\{0\}, \quad \forall t \in \mathbb{R} \cup\{\infty\}
$$

## Solutions of the eigenvalue equation

- The $q$-exponential function:

$$
\mathcal{E}_{q}(z)=\sum_{n=0}^{\infty} \frac{q^{n^{2} / 4}}{(q ; q)_{n}} z^{n}={ }_{1} \phi_{1}\left(0 ;-q^{1 / 2} ; q^{1 / 2},-q^{1 / 4} z\right)
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- Sequences $\psi^{ \pm}$, where

$$
\psi_{n}^{ \pm}:=( \pm \mathrm{i})^{n} q^{n / 2} \mathcal{E}_{q^{2}}\left( \pm \mathrm{i} x q^{n}\right)
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- By inspection of the asymptotic behavior of $\psi_{n}^{ \pm}$, as $n \rightarrow \pm \infty$, one gets:

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$$

- Hence, one expects there are non-trivial coefficients $a=a(x)$ and $b=b(x)$ such that

$$
a \psi^{+}+b \psi^{-} \in \ell^{2}(-\infty)
$$

## The $\ell^{2}$-solution

## Proposition

For all $x \in \mathbb{C} \backslash\{0\}$, the sequence

$$
\varphi(x):=\theta_{q}\left(-\mathrm{i} q^{-1 / 2} x\right) \psi^{(-)}(x)+\theta_{q}\left(\mathrm{i}^{-1 / 2} x\right) \psi^{(+)}(x)
$$

is the non-trivial solution of $\mathcal{A} \phi=x \phi$ which belongs to $\ell^{2}(\mathbb{Z})$. In addition, within the space $\ell^{2}(\mathbb{Z})$, this solution is given uniquely up to a multiplicative constant.

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Moreover,

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\varphi_{n}(x)=(-1 ; q)_{\infty} x^{n} q^{n(n-1) / 2}{ }_{0} \phi_{1}\left(-; 0 ; q^{2}, q^{-2 n+4} x^{-2}\right)
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and

$$
\|\varphi(x)\|_{\ell^{2}}^{2}=4 \frac{\left(q^{2} ; q^{2}\right)_{\infty}^{2}}{\left(q ; q^{2}\right)_{\infty}^{2}} \theta_{q^{2}}\left(-z^{2}\right)
$$

## The secular equation

## Theorem (secular equation)

For $t \in \mathbb{R} \cup\{\infty\}$, one has $\operatorname{spec}_{c}\left(A_{t}\right)=\{0\}$ and $\operatorname{spec}_{p}\left(A_{t}\right)$ coincides with the set of roots of the secular equation:

$$
x \theta_{q^{4}}\left(q^{2} x^{2}\right)+t \theta_{q^{4}}\left(x^{2}\right)=0, \quad \text { for } t \in \mathbb{R}
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and

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In addition, all eigenvalues of $A_{t}$ are simple.

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Corollary:

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\operatorname{spec}_{p}\left(A_{0}\right)=\left\{ \pm q^{2 n+1} \mid n \in \mathbb{Z}\right\} \quad \text { and } \quad \operatorname{spec}_{p}\left(A_{\infty}\right)=\left\{ \pm q^{2 n} \mid n \in \mathbb{Z}\right\}
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- How to solve the secular equation in general?
- Reparametrize $t=\Phi(s)$ and use nice properties of Jacobi's theta functions...


## Reparametrization $t=\Phi(s)$

- The function

$$
\Phi(s):=\mathrm{i} q^{1 / 2} \frac{\vartheta_{4}\left(\mathrm{is} \mid q^{2}\right)}{\vartheta_{1}\left(\mathrm{is} \mid q^{2}\right)}
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is real-valued, strictly decreasing on $(0,-2 \ln q)$, and maps $[0,-2 \ln q$ ) onto $\mathbb{R} \cup\{\infty\}$.

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- For the inverse function, one has

$$
\Phi^{-1}(t)=C(q) \int_{t}^{\infty} \frac{\mathrm{d} x}{\sqrt{\left(D(q)+x^{2}\right)\left(q+x^{2}\right)}}
$$

where

$$
C(q)=\frac{q^{1 / 2}}{\vartheta_{2}\left(0 \mid q^{2}\right) \vartheta_{3}\left(0 \mid q^{2}\right)}
$$

and

$$
D(q)=\frac{q \vartheta_{3}^{2}\left(0 \mid q^{2}\right)}{\vartheta_{2}^{2}\left(0 \mid q^{2}\right)}
$$

## Spectrum fully explicitly

- Using the reparametrization $t=\Phi(s)$, the secular equation simplifies to

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## Theorem

Let $t \in \mathbb{R} \cup\{\infty\}$, then

$$
\operatorname{spec}_{p}\left(A_{t}\right)=-e^{-s} q^{2 \mathbb{Z}} \cup e^{s} q^{2 \mathbb{Z}}
$$

where

$$
s=C(q) \int_{t}^{\infty} \frac{\mathrm{d} x}{\sqrt{\left(D(q)+x^{2}\right)\left(q+x^{2}\right)}}
$$

In addition, the family of corresponding eigenvectors $\left\{\varphi\left( \pm e^{s} q^{2 N}\right) \mid N \in \mathbb{Z}\right\}$, where

$$
\varphi_{n}(x)=(-1 ; q)_{\infty} x^{n} q^{n(n-1) / 2}{ }_{0} \phi_{1}\left(-; 0 ; q^{2}, q^{-2 n+4} x^{-2}\right)
$$

forms an orthogonal basis of $\ell^{2}(\mathbb{Z})$.

## (1) Introduction

2 Spectral resolution of A
(3) Spectral resolution of $B$

## The discrete Schrödinger operator $B$

The second Jacobi matrix $\mathcal{B}$ determines the unique operator

$$
B=U+U^{*}+\alpha V
$$

where $U$ is the forward shift operator and $V$ is the self-adjoint diagonal operator:

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U e_{n}=e_{n+1} \quad \text { and } \quad V e_{n}=q^{-n} e_{n}, \quad \forall n \in \mathbb{Z}
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## Proposition (essential spectrum)

The operator $B$ is self-adjoint and one has

$$
\sigma_{\mathrm{ess}}(B)=[-2,2] .
$$

## Solutions of the eigenvalue equation

- Using properties of the Hahn-Exton $q$-Bessel functions one verifies:

$$
f_{n}(z):=(-1)^{n} \alpha^{-n} q^{\frac{1}{2} n(n+1)}\left(z^{-1} \alpha^{-1} q^{n+1} ; q\right)_{\infty}{ }^{1} \phi_{1}\left(0 ; z^{-1} \alpha^{-1} q^{n+1} ; q, z \alpha^{-1} q^{n+1}\right)
$$

and

$$
g_{n}(z):=z^{-n}\left(z \alpha q^{1-n} ; q\right)_{\infty}{ }^{1} \phi_{1}\left(0 ; z \alpha q^{1-n} ; q, q z^{2}\right)
$$

are two solutions of the equation

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\mathcal{B} \psi=\left(z+z^{-1}\right) \psi .
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for all $\alpha, z \neq 0$.

## Solutions of the eigenvalue equation

- Using properties of the Hahn-Exton $q$-Bessel functions one verifies:

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- The solutions $f(z)$ and $g(z)$ are linearly independent iff $z \notin \alpha^{-1} q^{\mathbb{Z}} \cup\{0\}$ since

$$
W(f, g)=-z^{-1} \theta_{q}(\alpha z)
$$

## The point spectrum

Detailed asymptotic analysis of solutions $f$ and $g$ yields:

- If $0<|z|<1$ and $z \notin \alpha^{-1} q^{\mathbb{Z}} \cup\{0\}$, then

$$
f(z)\left\{\begin{array} { l } 
{ \in \ell ^ { 2 } ( + \infty ) } \\
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- If $|z|=1$ the asymptotic behavior of solutions is very different and, in the end, it implies that: For $\forall \alpha \in \mathbb{R}$ and $\forall x \in[-2,2]$, there is no non-trivial solution of $\mathcal{B} \psi=x \psi$ belonging to $\ell^{2}(\mathbb{Z})$.


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## Theorem (point spectrum)

If $\alpha \neq 0$, then

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\sigma(B) \backslash[-2,2]=\sigma_{p}(B)=\left\{\alpha^{-1} q^{m}+\alpha q^{-m} \mid m>\left\lfloor\log _{q}|\alpha|\right\rfloor\right\}
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In addition,

$$
\left\|\boldsymbol{v}_{m}\right\|_{\ell^{2}(\mathbb{Z})}=\frac{|\alpha|^{-m} q^{m(m+1) / 2}}{\sqrt{1-\alpha^{-2} q^{2 m}}}(q ; q)_{\infty}, \quad m>\left\lfloor\log _{q}|\alpha|\right\rfloor
$$

## The absolutely continuous part of the spectral measure

- Let us denote

$$
E_{k, l}(\cdot):=\left\langle e_{k}, E_{B}(\cdot) e_{l}\right\rangle, \quad k, I \in \mathbb{Z},
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where $E_{B}$ stands for the spectral measure of $B$.

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- To determine the spectral measure in the essential spectrum we use the formula

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E_{k, I}((a, b))=\lim _{\delta \rightarrow 0+\epsilon \rightarrow 0+} \lim _{2 \pi \mathrm{i}} \frac{1}{a+\delta} \int_{a-\delta}^{b-}\left(G_{k, I}(x+\mathrm{i} \epsilon)-G_{k, l}(x-\mathrm{i} \epsilon)\right) \mathrm{d} x
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where

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## Proposition

Let $\alpha \neq 0$ and $-2 \leq a<b \leq 2$. Then for any $k, l \in \mathbb{Z}$, it holds

$$
E_{k, l}([a, b])=\frac{1}{2 \pi} \int_{\phi_{b}}^{\phi_{a}} f_{l}\left(e^{\mathrm{i} \phi}\right) f_{k}\left(e^{\mathrm{i} \phi}\right)\left|\frac{\left(e^{2 \mathrm{i} \phi} ; q\right)_{\infty}}{\left(\alpha e^{\mathrm{i} \phi}, q \alpha^{-1} e^{-\mathrm{i} \phi} ; q\right)_{\infty}}\right|^{2} \mathrm{~d} \phi
$$

where $\phi_{a}=\arccos (a / 2)$ and $\phi_{b}=\arccos (b / 2)$. Consequently, $\sigma_{a c}(B)=[-2,2]$.

## Summary

## Theorem

If $\alpha \neq 0$, then

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\sigma_{e s s}(B)=\sigma_{a c}(B)=[-2,2],
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and

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In addition, for $\mathcal{M} \subset \mathbb{R}$ a Borel set, we have

$$
\begin{aligned}
E_{k, l}(\mathcal{M}) & =\frac{1}{2 \pi} \int_{2 \cos \phi \in[-2,2] \cap \mathcal{A}} f_{l}\left(e^{\mathrm{i} \phi}\right) f_{k}\left(e^{\mathrm{i} \phi}\right)\left|\frac{\left(e^{2 \mathrm{i} \phi} ; q\right)_{\infty}}{\left(\alpha e^{\mathrm{i} \phi}, q \alpha^{-1} e^{-\mathrm{i} \phi} ; q\right)_{\infty}}\right|^{2} \mathrm{~d} \phi \\
& +\frac{1}{(q ; q)_{\infty}^{2}} \sum_{\substack{m>\lfloor\log |\alpha|\rfloor \\
\alpha^{-1} q^{m}+\alpha q^{-m} \in \mathcal{M}}}\left(1-\alpha^{-2} q^{2 m}\right) \alpha^{2 m} q^{-m(m+1)} f_{l}\left(\alpha^{-1} q^{m}\right) f_{k}\left(\alpha^{-1} q^{m}\right)
\end{aligned}
$$

## One consequence for special functions

- Recall the Hanhn-Exton (or third Jackson's) $q$-Bessel function is defined as

$$
J_{\nu}(z ; q)=z^{\nu} \frac{\left(q^{\nu+1} ; q\right)_{\infty}}{(q ; q)_{\infty}} 1 \phi_{1}\left(0 ; q^{\nu+1} ; q, q z^{2}\right)
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- Elements of the eigenvectors $\boldsymbol{v}_{m}$ are expressible in terms of Hanhn-Exton $q$-Bessel function and the formula for the norm of the eigenvectors yields

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\sum_{n \in \mathbb{Z}} J_{n}^{2}(z ; q)=\frac{1}{1-z^{2}}, \quad|z|<1
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- This formula seems to be new (Really?) and it generalizes the well-known summation formula for the Bessel functions of the first kind:

$$
\sum_{n \in \mathbb{Z}} J_{n}^{2}(z)=1, \quad|z|<1
$$

## Thank you!

